

Landé g-factor Derivation.
R. Corn -- March 2020.

In the weak field limit, we assume that the magnetic dipole moment due to the electron in an atom is proportional to the total angular momentum \mathbf{J} :

$$\boldsymbol{\mu}_J = -g_J \frac{m_B}{\hbar} \mathbf{J}$$

where m_B is the Bohr magneton ($\frac{e\hbar}{2m_e}$) and g_J is the Landé g factor for \mathbf{J} . To find a value of g_J , we relate \mathbf{J} to the known values of g_L and g_S and assume the following:

$$\begin{aligned} \mathbf{J} &= \mathbf{L} + \mathbf{S} & \boldsymbol{\mu}_J &= \boldsymbol{\mu}_L + \boldsymbol{\mu}_S \\ \boldsymbol{\mu}_L &= -g_L \frac{m_B}{\hbar} \mathbf{L} & \boldsymbol{\mu}_S &= -g_S \frac{m_B}{\hbar} \mathbf{S} \\ \boldsymbol{\mu}_J &= -g_J \frac{m_B}{\hbar} \mathbf{J} = -g_L \frac{m_B}{\hbar} \mathbf{L} + -g_S \frac{m_B}{\hbar} \mathbf{S} \\ g_J \mathbf{J} &= g_L \mathbf{L} + g_S \mathbf{S} \end{aligned}$$

We then take the dot product with \mathbf{J} and get:

$$\begin{aligned} g_J \mathbf{J} \cdot \mathbf{J} &= g_L \mathbf{L} \cdot \mathbf{J} + g_S \mathbf{S} \cdot \mathbf{J} \\ g_J J^2 &= g_L (L^2 + \mathbf{L} \cdot \mathbf{S}) + g_S (S^2 + \mathbf{L} \cdot \mathbf{S}) \\ g_J J^2 &= g_L \left(L^2 + 0.5 (J^2 - L^2 - S^2) \right) + g_S \left(S^2 + 0.5 (J^2 - L^2 - S^2) \right) \\ J^2 &= \hbar^2 J(J+1) & L^2 &= \hbar^2 L(L+1) & S^2 &= \hbar^2 S(S+1) \\ g_J &= \frac{g_L (J(J+1) + L(L+1) - S(S+1)) + g_S (J(J+1) - L(L+1) + S(S+1))}{2J(J+1)} \end{aligned}$$

For the case:

$$g_L = 1$$

$$g_S = 2$$

We get:

$$g_J = \frac{3}{2} + \frac{S(S+1) - L(L+1)}{2J(J+1)}$$

And so, for the weak field limit:

$$W_z = -\boldsymbol{\mu}_J \cdot \mathbf{B} = g_J \omega_0 J_z$$

$$\omega_0 = -\frac{m_B}{\hbar} B_0 = -\frac{eB_0}{2m_e}$$

where ω_0 is the Larmor angular frequency.

$g_J = 4/3$ for 2p($j=3/2$) state ($L=1, S=1/2, J=3/2$)

$g_J = 2/3$ for 2p($j=1/2$) state ($L=1, S=1/2, J=1/2$)

Et voilà! These are eqns. 11-13 in the CTcompD12.pdf handout.