

## CHAPTER 6: DIFFUSION IN SOLIDS

Gear from case-hardened steel (C diffusion)



### ISSUES TO ADDRESS...

- How does diffusion occur?
- Why is it an important part of processing?
- How can the rate of diffusion be predicted for some simple cases?
- How does diffusion depend on structure and temperature?

## Diffusion- Steady and Non-Steady State

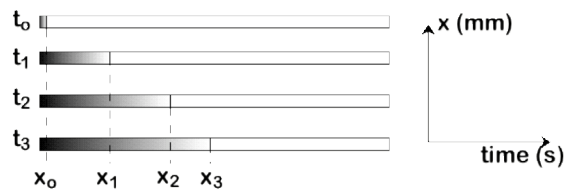
Diffusion - Mass transport by atomic motion

### Mechanisms

- Gases & Liquids – random (Brownian) motion
- Solids – vacancy diffusion or interstitial diffusion

## Simple Diffusion

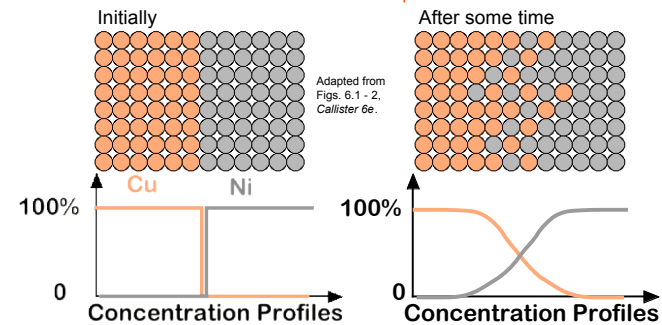
- Glass tube filled with water.
- At time  $t = 0$ , add some drops of ink to one end of the tube.
- Measure the diffusion distance,  $x$ , over some time.
- Compare the results with theory.



## Inter-diffusion

- **Interdiffusion:** In alloys, atoms tend to migrate from regions of large concentration.

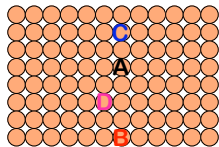
This is a diffusion couple.



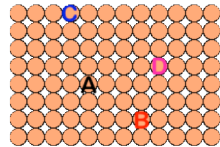
## Self-diffusion

- **Self-diffusion:** In an elemental solid, atoms also migrate.

Label some atoms



After some time



## Substitution-diffusion: vacancies and interstitials

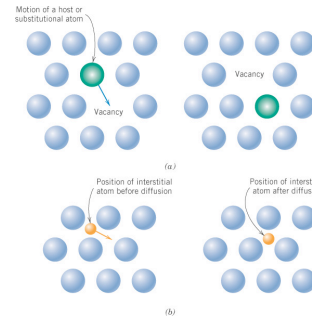
- applies to substitutional impurities
- atoms exchange with vacancies
- **rate** depends on (1) **number of vacancies**;  
(2) **activation energy to exchange.**

**Number** (or concentration\*)  
**of Vacancies** at T

$$c_v = \frac{n_v}{N} = e^{-\frac{\Delta E}{k_B T}}$$

\*  $k_B T$  gives eV

\* see web handout for derivation.

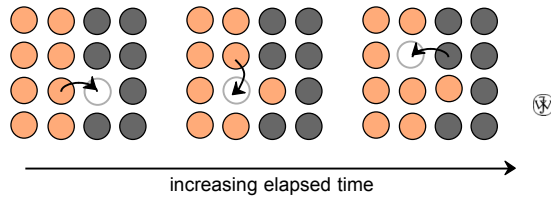


$\Delta E$  is an **activation energy**  
for a particular process  
(in J/mol, cal/mol, eV/atom).

## Substitution-diffusion

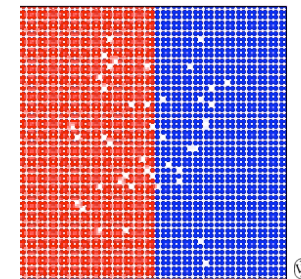
**Vacancy Diffusion:**

- applies to substitutional impurities
- atoms exchange with vacancies
- **rate** depends on (1) **number of vacancies**;  
(2) **activation energy to exchange.**



## Inter-diffusion across Interfaces

- Rate of substitutional diffusion depends on:
  - vacancy concentration
  - frequency of jumping.

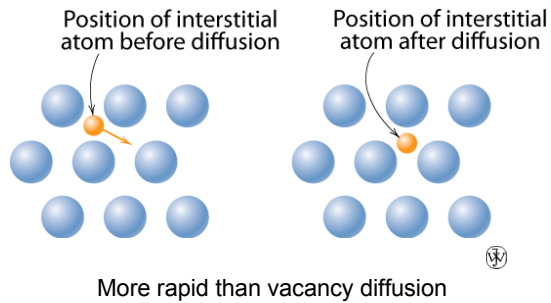


(Courtesy P.M. Anderson)

Why should interstitial diffusion be faster than by vacancy mode of diffusion?

## Diffusion Mechanisms

- **Interstitial diffusion** – smaller atoms diffuse between atoms.

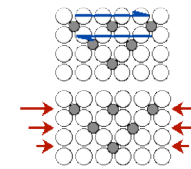


## Processing using Diffusion

- **Case Hardening:**
  - Diffuse carbon atoms into the host iron atoms at the surface.
  - Example of interstitial diffusion is a case hardened gear.
- **Result:** The "Case" is
  - **hard to deform:** C atoms "lock" planes from **shearing**.
  - **hard to crack:** C atoms put the surface in compression.



Fig. 6.0, Callister 6e.  
(courtesy of Surface Div., Midland-Ross.)



## Processing using Diffusion

- **Doping Silicon with P** for n-type semiconductors:
- **Process:**

1. Deposit P rich layers on surface.



2. Heat it.
3. Result: Doped semiconductor regions.

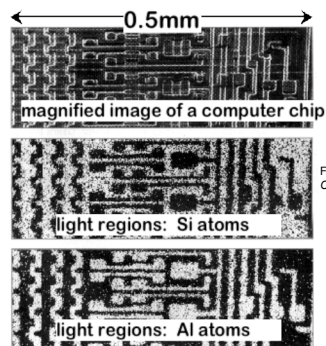
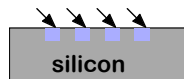
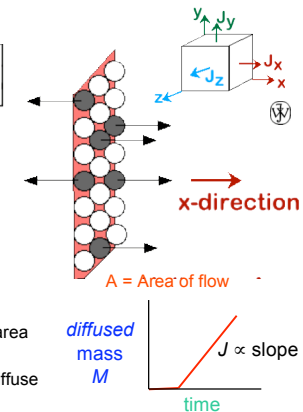


Fig. 18.0,  
Callister 6e.

## Modeling rate of diffusion: flux

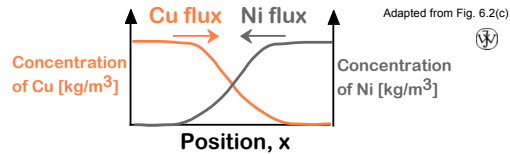
- **Flux:**

$$J = \frac{1}{A} \frac{dM}{dt} \Rightarrow \left[ \frac{\text{kg}}{\text{m}^2\text{s}} \right] \text{ or } \left[ \frac{\text{atoms}}{\text{m}^2\text{s}} \right]$$
- **Directional Quantity**
- **Flux can be measured for:**
  - vacancies
  - host (A) atoms
  - impurity (B) atoms
- **Empirically determined:**
  - Make thin membrane of known surface area
  - Impose concentration gradient
  - Measure how fast atoms or molecules diffuse through the membrane



## Steady-state Diffusion: $J \sim$ gradient of $c$

- Concentration Profile,  $C(x)$ : [ $\text{kg}/\text{m}^3$ ]



- Fick's First Law:  $D$  is a constant!

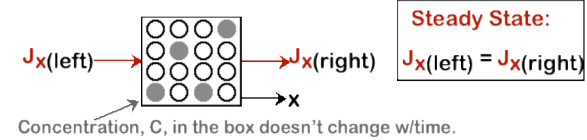
$$\text{flux in } x\text{-dir. [kg/m}^2\text{-s]} \rightarrow J_x = -D \frac{dC}{dx}$$

$D$  ← Diffusion coefficient [ $\text{m}^2/\text{s}$ ]  
 $\frac{dC}{dx}$  ← concentration gradient [ $\text{kg}/\text{m}^3$ ]

- The steeper the concentration profile, the greater the flux!

## Steady-State Diffusion

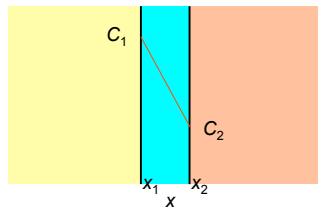
- Steady State: concentration profile not changing with time.



- Apply Fick's First Law:  $J_x = -D \frac{dC}{dx}$
- If  $J_x(\text{left}) = J_x(\text{right})$ , then  $\left(\frac{dC}{dx}\right)_{\text{left}} = \left(\frac{dC}{dx}\right)_{\text{right}}$
- Result: the slope,  $dC/dx$ , must be constant (i.e., slope doesn't vary with position)!

## Steady-State Diffusion

Rate of diffusion independent of time  $J \sim \frac{dC}{dx}$



Fick's first law of diffusion

$$J = -D \frac{dC}{dx}$$

$D$  = diffusion coefficient

if linear  $\frac{dC}{dx} \approx \frac{\Delta C}{\Delta x} = \frac{C_2 - C_1}{x_2 - x_1}$

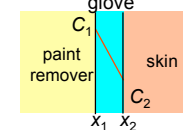
## Example: Chemical Protection Clothing

- Methylene chloride is a common ingredient of paint removers. Besides being an irritant, it also may be absorbed through skin. When using, protective gloves should be worn.

- If butyl rubber gloves (0.04 cm thick) are used, what is the diffusive flux of methylene chloride through the glove?

- Data:

- $D$  in butyl rubber:  $D = 110 \times 10^{-8} \text{ cm}^2/\text{s}$
- surface concentrations:  $C_1 = 0.44 \text{ g}/\text{cm}^3$   $C_2 = 0.02 \text{ g}/\text{cm}^3$
- Diffusion distance:  $x_2 - x_1 = 0.04 \text{ cm}$



$$J_x = -D \frac{C_2 - C_1}{x_2 - x_1} = 1.16 \times 10^{-5} \frac{\text{g}}{\text{cm}^2\text{s}}$$

### Example: C Diffusion in steel plate

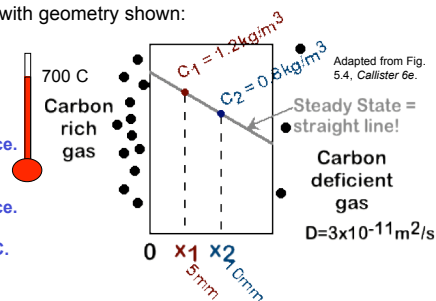
- Steel plate at 700°C with geometry shown:

**Knowns:**

$C_1 = 1.2 \text{ kg/m}^3$  at 5mm  
( $5 \times 10^{-3} \text{ m}$ ) below surface.

$C_2 = 0.8 \text{ kg/m}^3$  at 10mm  
( $1 \times 10^{-2} \text{ m}$ ) below surface.

$D = 3 \times 10^{-11} \text{ m}^2/\text{s}$  at 700 C.



Adapted from Fig. 5.4, Callister 6e.

- Q: In steady-state, how much carbon transfers from the rich to the deficient side?

$$J = -D \frac{C_2 - C_1}{x_2 - x_1} = 2.4 \times 10^{-9} \frac{\text{kg}}{\text{m}^2 \cdot \text{s}}$$

### Example: Diffusion of radioactive atoms

- Surface of Ni plate at 1000°C contains 50% Ni<sup>63</sup> (radioactive) and 50% Ni (non-radioactive).
- 4 microns below surface Ni<sup>63</sup> / Ni = 48:52
- Lattice constant of Ni at 1000 C is 0.360 nm.
- Experiment shows that self-diffusion of Ni is  $1.6 \times 10^{-9} \text{ cm}^2/\text{sec}$

What is the flux of Ni<sup>63</sup> atoms through a plane 2 μm below surface?

$$C_1 = \frac{(4 \text{ Ni / cell})(0.5 \text{ Ni}^{63} / \text{Ni})}{(0.36 \times 10^{-9} \text{ m})^3 / \text{cell}} = 42.87 \times 10^{27} \text{ Ni}^{63} / \text{m}^3$$

$$C_2 = \frac{(4 \text{ Ni / cell})(0.48 \text{ Ni}^{63} / \text{Ni})}{(0.36 \times 10^{-9} \text{ m})^3 / \text{cell}} = 41.15 \times 10^{27} \text{ Ni}^{63} / \text{m}^3$$

$$J = -D \frac{C_2 - C_1}{x_2 - x_1} = -(1.6 \times 10^{-13} \text{ m}^2 / \text{sec}) \frac{(41.15 - 42.87) \times 10^{27} \text{ Ni}^{63} / \text{m}^3}{(4 - 0) \times 10^{-6} \text{ m}} = -0.69 \times 10^{20} \text{ Ni}^{63} / \text{m}^2 \cdot \text{s}$$

How many Ni<sup>63</sup> atoms/second through cell?

$$J \cdot (0.36 \text{ nm})^2 = 9 \text{ Ni}^{63} / \text{s}$$

### Where can we use Fick's Law?

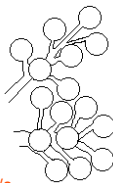
Fick's law is commonly used to model transport processes in

- foods,
- clothing,
- biopolymers,
- pharmaceuticals,
- porous soils,
- semiconductor doping process, etc.

**Example**

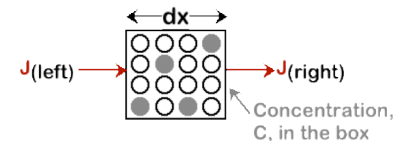
The total membrane surface area in the lungs (alveoli) may be on the order of **100 square meters** and have a thickness of less than a **millionth of a meter**, so it is a very effective **gas-exchange interface**.

CO<sub>2</sub> in air has D~16 mm<sup>2</sup>/s, and, in water, D~ 0.0016 mm<sup>2</sup>/s



### Non-Steady-State Diffusion

- Concentration profile, C(x), changes w/ time.



- To conserve matter:

$$\frac{J(\text{right}) - J(\text{left})}{dx} = -\frac{dC}{dt}$$

$$\frac{dJ}{dx} = -\frac{dC}{dt}$$

- Fick's First Law:

$$J = -D \frac{dC}{dx} \text{ or}$$

$$\frac{dJ}{dx} = -D \frac{d^2C}{dx^2} \text{ (if D does not vary with x)}$$

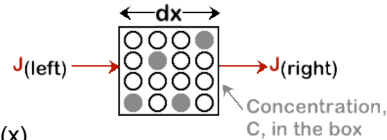
equate

- Governing Eqn.:

$$\frac{dC}{dt} = D \frac{d^2C}{dx^2}$$

## Non-Steady-State Diffusion: another look

- Concentration profile,  $C(x)$ , changes w/ time.



- Rate of accumulation  $C(x)$

$$\frac{\partial C}{\partial t} dx = J_x - J_{x+dx} \rightarrow \frac{\partial C}{\partial t} dx = J_x - (J_x + \frac{\partial J_x}{\partial x} dx) = -\frac{\partial J_x}{\partial x} dx$$

- Using Fick's Law:  $\frac{\partial C}{\partial t} = -\frac{\partial J_x}{\partial x} = -\frac{\partial}{\partial x} \left( -D \frac{\partial C}{\partial x} \right)$  Fick's 2nd Law

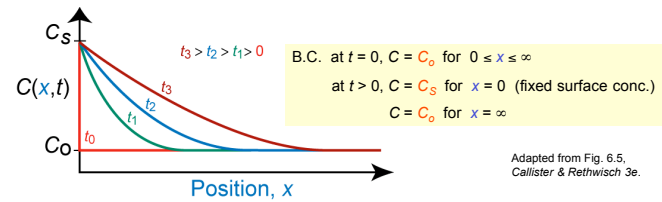
- If  $D$  is constant: Fick's Second "Law"  $\frac{\partial C}{\partial t} = \frac{\partial}{\partial x} \left( D \frac{\partial C}{\partial x} \right) \approx \left( D \frac{\partial^2 C}{\partial x^2} \right)$

## Non-Steady-State Diffusion: $C = c(x,t)$

concentration of diffusing species is a function of both time and position

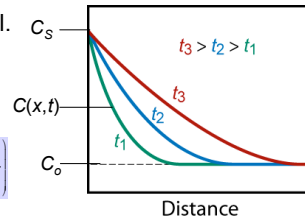
$$\text{Fick's Second "Law"} \quad \frac{\partial C}{\partial t} \approx \left( D \frac{\partial^2 C}{\partial x^2} \right)$$

- Copper diffuses into a bar of aluminum.



## Non-Steady-State Diffusion

- Cu diffuses into a bar of Al.



$$\text{Fick's Second "Law": } \frac{\partial C}{\partial t} \approx \left( D \frac{\partial^2 C}{\partial x^2} \right)$$

$$\text{Solution: } \frac{C(x,t) - C_o}{C_s - C_o} = 1 - \text{erf} \left( \frac{x}{2\sqrt{Dt}} \right)$$

"error function" Values calibrated in Table 6.1

$$\text{erf}(z) = \frac{2}{\sqrt{\pi}} \int_0^z e^{-y^2} dy$$

## Example: Non-Steady-State Diffusion

FCC iron-carbon alloy initially containing 0.20 wt% C is carburized at an elevated temperature and in an atmosphere that gives a surface C content at 1.0 wt%.

If after 49.5 h the concentration of carbon is 0.35 wt% at a position 4.0 mm below the surface, what temperature was treatment done?

Solution

$$\frac{C(x,t) - C_o}{C_s - C_o} = \frac{0.35 - 0.20}{1.0 - 0.20} = 1 - \text{erf} \left( \frac{x}{2\sqrt{Dt}} \right) = 1 - \text{erf}(z) \quad \therefore \text{erf}(z) = 0.8125$$

Using Table 6.1 find  $z$  where  $\text{erf}(z) = 0.8125$ . Use interpolation.

$z$	$\text{erf}(z)$
0.90	0.7970
$z$	0.8125
0.95	0.8209

$$\frac{z - 0.90}{0.95 - 0.90} = \frac{0.8125 - 0.7970}{0.8209 - 0.7970} \quad \text{So, } z = 0.93$$

$$\text{Now solve for } D \quad z = \frac{x}{2\sqrt{Dt}} \Rightarrow D = \frac{x^2}{4z^2t}$$

$$\therefore D = \frac{x^2}{4z^2t} = \frac{(4 \times 10^{-3} \text{ m})^2}{(4)(0.93)^2 (49.5 \text{ h}) (3600 \text{ s})} = 2.6 \times 10^{-11} \text{ m}^2/\text{s}$$

**Solution (cont.):**

- To solve for the temperature at which  $D$  has the above value, we use a rearranged form of Equation (6.9a):  

$$D = D_0 \exp(-Q_d/RT)$$

$$T = \frac{Q_d}{R(\ln D_0 - \ln D)}$$
- From Table 6.2, for diffusion of C in FCC Fe  
 $D_0 = 2.3 \times 10^{-5} \text{ m}^2/\text{s}$      $Q_d = 148,000 \text{ J/mol}$      $D = 2.6 \times 10^{-11} \text{ m}^2/\text{s}$

$$T = \frac{148,000 \text{ J/mol}}{(8.314 \text{ J/mol}\cdot\text{K})(\ln 2.3 \times 10^{-5} \text{ m}^2/\text{s} - \ln 2.6 \times 10^{-11} \text{ m}^2/\text{s})}$$

$T = 1300 \text{ K} = 1027^\circ\text{C}$

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### Example: Processing

- Copper diffuses into a bar of aluminum.
- 10 hours processed at 600 C gives desired  $C(x)$ .
- How many hours needed to get the same  $C(x)$  at 500 C?

Key point 1:  $C(x, t_{500C}) = C(x, t_{600C})$ .  
 Key point 2: Both cases have the same  $C_0$  and  $C_s$ .

Result:  $Dt$  should be held constant.

$$\frac{C(x,t) - C_0}{C_s - C_0} = 1 - \text{erf}\left(\frac{x}{2\sqrt{Dt}}\right) \rightarrow (Dt)_{500^\circ\text{C}} = (Dt)_{600^\circ\text{C}}$$

$5.3 \times 10^{-13} \text{ m}^2/\text{s} \rightarrow (Dt)_{600} = 10 \text{ hrs}$   
 $4.8 \times 10^{-14} \text{ m}^2/\text{s} \rightarrow (Dt)_{500} = 110 \text{ hr}$

Answer:  $t_{500} = \frac{(Dt)_{600}}{D_{500}} = 110 \text{ hr}$

**Note**  
 $D(T)$  are T dependent!  
 Values of  $D$  are provided.

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### Diffusion Analysis

- The experiment: we recorded combinations of  $t$  and  $x$  that kept  $C$  constant.

$$\frac{C(x_i, t_i) - C_0}{C_s - C_0} = 1 - \text{erf}\left(\frac{x_i}{2\sqrt{Dt_i}}\right) = (\text{constant here})$$

- Diffusion depth given by:  $x_i \propto \sqrt{Dt_i}$

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### Data from Diffusion Analysis

- Experimental result:  $x \sim t^{0.58}$
- Theory predicts  $x \sim t^{0.50}$  from  $x_i \propto \sqrt{Dt_i}$
- Close agreement.

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## Diffusion and Temperature

- Diffusivity increases with T exponentially (so does Vacancy conc.).

$$D = D_0 \exp\left(-\frac{Q_d}{RT}\right)$$

$D$  = diffusion coefficient [ $m^2/s$ ]

$D_0$  = pre-exponential [ $m^2/s$ ] (see Table 6.2)

$Q_d$  = activation energy [J/mol or eV/atom]

$R$  = gas constant [8.314 J/mol-K]

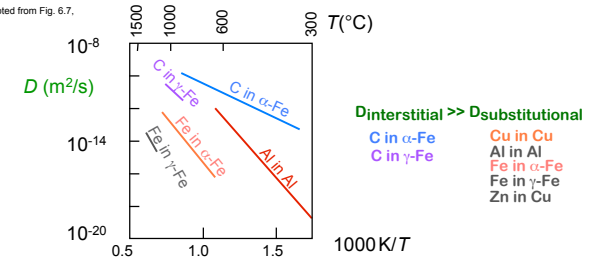
$T$  = absolute temperature [K]

Note:  $\ln D = \ln D_0 - \frac{Q_d}{R} \left(\frac{1}{T}\right)$        $\log D = \log D_0 - \frac{Q_d}{2.3R} \left(\frac{1}{T}\right)$

## Diffusion and Temperature

- Experimental Data:

Adapted from Fig. 6.7.



Adapted from Fig. 6.7, Callister & Rethwisch 3e.  
(Data for Fig. 6.7 from E.A. Brandes and G.B. Brook (Ed.) *Smithells Metals Reference Book*, 7th ed., Butterworth-Heinemann, Oxford, 1992.)

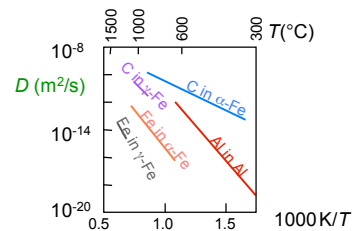
## Example: Comparing Diffuse in Fe

Is C in fcc Fe diffusing faster than C in bcc Fe?

(Table 6.2)

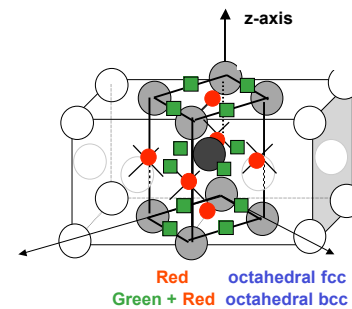
- fcc-Fe:  $D_0 = 2.3 \times 10^{-5} (m^2/2)$   
 $Q = 1.53 \text{ eV/atom}$   
 $T = 900 \text{ C}$   $D = 5.9 \times 10^{-12} (m^2/2)$

- bcc-Fe:  $D_0 = 6.2 \times 10^{-7} (m^2/2)$   
 $Q = 0.83 \text{ eV/atom}$   
 $T = 900 \text{ C}$   $D = 1.7 \times 10^{-10} (m^2/2)$



- FCC Fe has both higher activation energy  $Q$  and  $D_0$  (holes larger in FCC).
- BCC and FCC phase exist over limited range of  $T$  (at varying %C).
- Hence, at same  $T$ , BCC diffuses faster due to lower  $Q$ .
- Cannot always have the phase that you want at the %C and  $T$  you want! which is why this is all important.

## Connecting Holes, Diffusion, and Stress



FCC represented as a BCT cell, with relevant holes.



### Other Types of Diffusion (Beside Atomic)

- Flux is general concept: e.g. charges, phonons,...

Charge Flux –  $j_q = \frac{1}{A} \frac{dq}{dt} = \frac{1}{A} \frac{d(Ne)}{dt}$        $e = \text{electric chg.}$   
 $N = \text{net \# e- cross A}$

Defining conductivity  $\sigma$  (a material property)  $j = -\sigma \frac{dV}{dx}$

Ohm's Law  $j_q = \frac{I}{A} = \frac{1}{AR} \frac{dV}{dx} = \sigma \frac{V}{\ell}$

Solution: Fick's 2nd Law

$$\frac{V_x - V_s}{V_0 - V_s} = \text{erf}\left(\frac{x}{2\sqrt{\sigma t}}\right)$$

Heat Flux – (by phonons)  $Q = \frac{1}{A} \frac{dH}{dt} = \frac{1}{A} \frac{d(N\varepsilon)}{dt}$

$N = \# \text{ of phonons with avg. energy } \varepsilon$

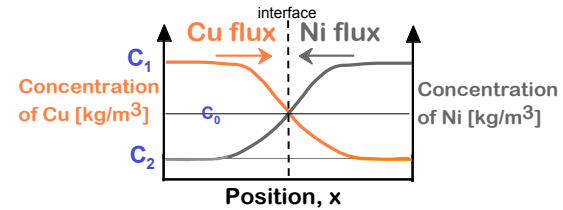
Defining thermal conductivity  $\kappa$  (a material property)  $Q = -\kappa \frac{dT}{dx}$

Solution: Fick's 2nd Law

$$\frac{T_x - T_s}{T_0 - T_s} = \text{erf}\left(\frac{x}{2\sqrt{\kappa t}}\right)$$

Or w/ Thermal Diffusivity:  $h = \frac{\kappa}{c_v \rho}$        $Q = c_v \rho T$

### Inter-diffusion (diffusion couples)

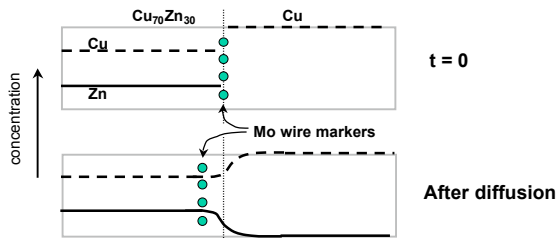


Assuming  $D_A = D_B$ :  $\frac{C_1^x - C_0}{C_1 - C_0} = \text{erf}\left(\frac{x}{2\sqrt{Dt}}\right)$        $C_0 = (C_1 + C_2)/2$

For  $C_2^x$  Curve is symmetric about  $C_0$ ,  $\text{erf}(-z) = -\text{erf}(z)$

### Kirkendall Effect: What is $D_A > D_B$ ?

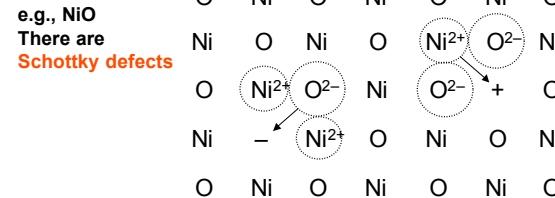
- Kirkendall studied Mo markers in Cu-brass (i.e., fcc  $\text{Cu}_{70}\text{Zn}_{30}$ ).
- Symmetry is lost: Zn atoms move more readily in one direction (to the right) than Cu atoms move in the other (to the left).



- When diffusion is asymmetric, interface moves away markers, i.e., there is a net flow of atoms to the right past the markers.
- Analyzing movement of markers determines  $D_{\text{Zn}}$  and  $D_{\text{Cu}}$ .
- Kirkendall effect driven by vacancies, effective for  $T > 0.5 T_{\text{melt}}$ .

### Diffusion in Compounds: Ionic Conductors

- Unlike diffusion in metals, diffusion in compounds involves second-neighbor migration.
- Since the activation energies are high, the  $D$ 's are low unless vacancies are present from non-stoichiometric ratios of atoms.



- The two vacancies cannot accept neighbors because they have wrong charge, and ion diffusion needs 2nd neighbors with high barriers (activation energies).

## Diffusion in Compounds: Ionic Conductors

- D's in an ionic compound are seldom comparable because of **size, change and/or structural differences**.
- Two sources of conduction: **ion diffusion** and via **e<sup>-</sup> hopping from ions of variable valency**, e.g., Fe<sup>2+</sup> to Fe<sup>3+</sup>, in applied electric field.

e.g., **ionic**

- In **NaCl** at 1000 K,  $D_{Na^+} \sim 5D_{Cl^-}$ , whereas at 825 K  $D_{Na^+} \sim 50D_{Cl^-}$ !
- This is primarily due to size  $r_{Na^+} = 1 \text{ \AA}$  vs  $r_{Cl^-} = 1.8 \text{ \AA}$ .

e.g., **oxides**

- In **uranium oxide**, U<sup>4+</sup>(O<sup>2-</sup>)<sub>2</sub>, at 1000 K (extrapolated),  $D_O \sim 10^7 D_U$ .
- This is mostly due to charge, i.e. more energy to activate 4+ U ion.
- Also, UO is not stoichiometric, having U<sup>3+</sup> ions to give UO<sub>2-x</sub>, so that the anion vacancies significantly increase O<sup>2-</sup> mobility.

e.g., **solid-solutions of oxides (leads to defects, e.g., vacancies)**

- If **Fe<sub>1-x</sub>O** (x=2.5-4% at 1500 K, 3Fe<sup>2+</sup> → 2Fe<sup>3+</sup> + vac.) is dissolved in MgO under reducing conditions, then Mg<sup>2+</sup> diffusion increases.
- If **MgF<sub>2</sub> is dissolved in LiF** (2Li<sup>+</sup> → Mg<sup>2+</sup> + vac.), then Li<sup>+</sup> diffusion increases. All due to additional vacancies.

## Ionic Conduction: related to fuel cells

- Molten salts and aqueous electrolytes conduct charge when placed in electric field, +q and -q move in opposite directions.

- The same occurs in solids although at much slower rate.

- Each ion has **charge of Ze** ( $e = 1.6 \times 10^{-19} \text{ amp*sec}$ ),

so ion movement induces ionic conduction

- Conductivity**  $\sigma = n\mu Ze$  is related to **mobility,  $\mu$** , which is related to D via the Einstein equations:  $\mu = ZeD / k_B T$

- Hence

$$\sigma_{ionic} = \frac{nZ^2 e^2}{k_B T} D = \frac{nZ^2 e^2}{k_B T} D_0 e^{-Q/RT}$$

$$\log_{10} \sigma_{ionic} \sim \ln \left( \frac{nZ^2 e^2}{k_B T} D_0 \right) - \frac{Q}{2.3RT}$$

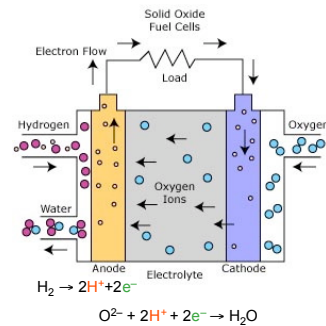
So, electrical conduction can be used to determine diffusion data in ionic solids.

e.g., **What conductivity results by Ca<sup>2+</sup> diffusion in CaO at 2000 K?**

- CaO has NaCl structure with  $a = 4.81 \text{ \AA}$ , with  $D(2000 \text{ K}) \sim 10^{-14} \text{ m}^2/\text{s}$ , and  $Z=2$ .

$$n_{Ca^{2+}} = \frac{4}{\text{cell}} \frac{\text{cell}}{(4.81 \times 10^{-10} \text{ m})^3} = 3.59 \times 10^{28} / \text{m}^3 \quad \sigma = \frac{nZ^2 e^2}{k_B T} D \sim \frac{1.3 \times 10^{-5}}{\text{ohm-cm}}$$

## Example: solid-oxide fuel cell (SOFC)



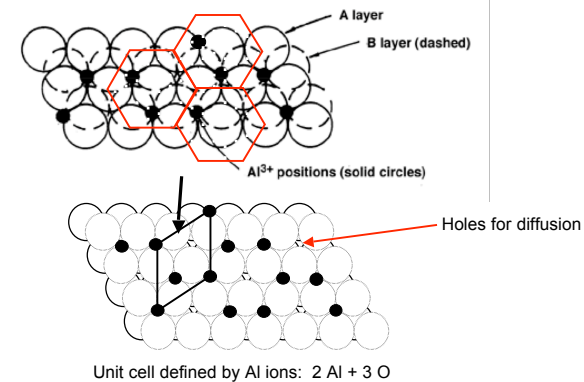
- SOFC is made up of four layers, three of which are ceramics (hence the name).
- A single cell consisting of these four layers stacked together is only a few mm thick.
- Need to stack many, many together to have larger DC current.



Image: <http://www.ip3.unipg.it/FuelCells/en/hfcs.asp>

S. Haile's SOFC (2004 best): Thin-film of Sm-doped Ceria electrolyte (CeO<sub>2</sub>, i.e. Sm<sub>x</sub>Ce<sub>1-x</sub>O<sub>2-x/2</sub>) and BSCF cathode (Perovskite Ba<sub>0.2</sub>Si<sub>0.5</sub>Co<sub>0.8</sub>Fe<sub>0.2</sub>O<sub>3-d</sub>) show high power densities – over 1 W/cm<sup>2</sup> at 600 C – with humidified H<sub>2</sub> as the fuel and air at the cathode.

## Ceramic Compounds: Al<sub>2</sub>O<sub>3</sub>



### Summary: Structure and Diffusion

Diffusion **FASTER** for...

- open crystal structures
- lower melting T materials
- materials w/secondary bonding
- smaller diffusing atoms
- cations
- lower density materials

Diffusion **SLOWER** for...

- close-packed structures
- higher melting T materials
- materials w/covalent bonding
- larger diffusing atoms
- anions
- higher density materials